

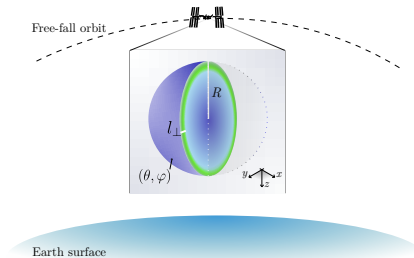
Quantum statistics and BKT transition of a shell-shaped superfluid

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Young Seminars SIFS



In collaboration with
A. Pelster, F. Cinti, L. Salasnich.

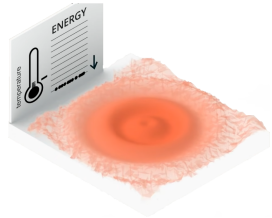
This presentation will be on
www.andreatononi.com

Outline

- ▷ Introduction and motivation
- ▷ Bose-Einstein condensation in shell-shaped condensates
 - ▶ Spherical case
 - ▶ Ellipsoidal case
- ▷ Superfluidity in 2D: BKT transition
 - ▶ Spherical superfluid
- ▷ Conclusions

Quantum gases

Bose & Einstein [1924-25]:
predicted the phenomenon of
Bose-Einstein condensation

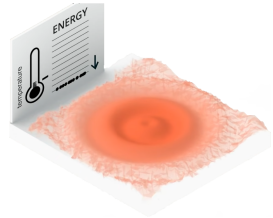


macroscopic occupation of the same lowest-energy
single-particle state of a many-body system of identical
bosons

(occupation numbers controlled by T , and by n)

Quantum gases

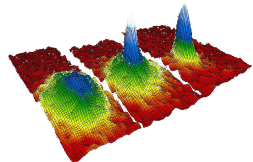
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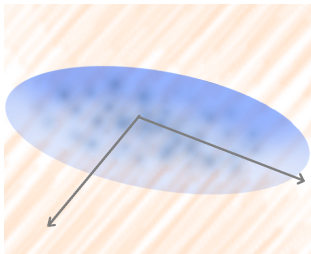
(occupation numbers controlled by T , and by n)

Cornell & Wieman, Ketterle, [1995]:
observation of BEC by trapping and
cooling alkali-metal atoms

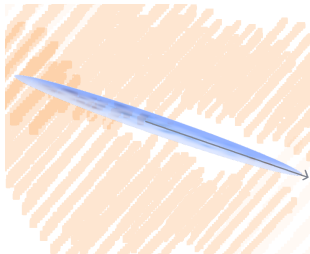


Low-dimensional quantum gases

(2D)

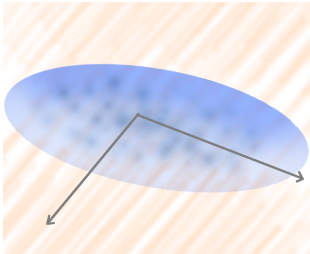


(1D)

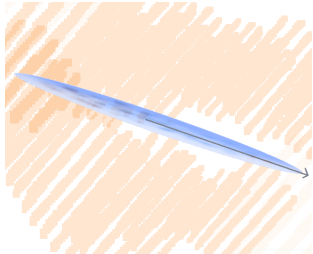


Low-dimensional quantum gases

(2D)



(1D)



Quantum many-body physics has been studied **consistently** only in “*flat*” low-dimensional configurations

What about *curved* geometries?

Shell-shaped quantum gases (rf-induced adiabatic potentials)

Theoretical proposal of [Zobay, Garraway, PRL **86**, 1195 (2001)]:
confine the atoms with $B_0(\vec{r})$, and $B_{rf}(\vec{r}, t)$, yielding

$$U(\vec{r}) = M_F \sqrt{\left[\sum_i \frac{m}{2} \omega_i^2 x_i^2 - \hbar \Delta \right]^2 + (\hbar \Omega)^2}$$

ω_i : frequencies of the bare harmonic trap

Δ : detuning from the resonant frequency

Ω : Rabi frequency between coupled levels

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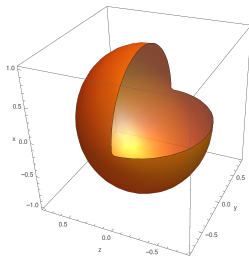
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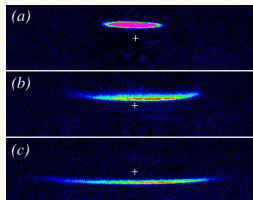
Minimum of $U(\vec{r})$ for

$$\omega_x^2 x^2 + \omega_y^2 y^2 + \omega_z^2 z^2 = \frac{2\hbar \Delta}{m}.$$



Shell-shaped quantum gases

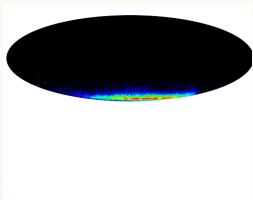
On Earth...



[Colombe *et al.*, EPL **67**, 593 (2004)]

Shell-shaped quantum gases

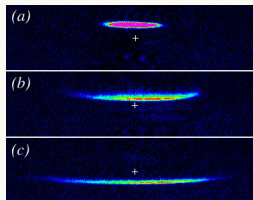
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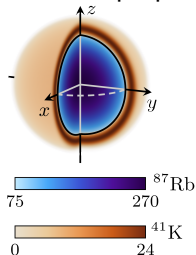
Shell-shaped quantum gases, in microgravity

On Earth...



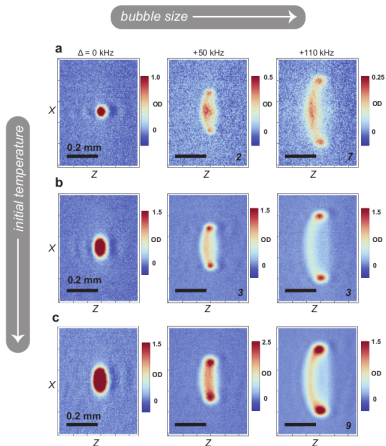
[Colombe *et al.*, EPL 67, 593 (2004)]

...and other proposals:



[Wolf *et al.*, arXiv:2110.15247]

...in microgravity:



[Carollo *et al.*, arXiv:2108.05880]

(experiments on CAL, on ISS)

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Bose-Einstein condensation on the surface of a sphere

Noninteracting case, single particle on a sphere of radius R :

$$\frac{\hat{L}^2}{2mR^2} \mathcal{Y}_{l,m_l}(\theta, \varphi) = \epsilon_l \mathcal{Y}_{l,m_l}(\theta, \varphi),$$

with $\epsilon_l = \frac{\hbar^2}{2mR^2} l(l+1)$ and $m_l = -l, \dots, +l$.

$$\text{Particle number: } N = N_0 + \sum_{l=1}^{+\infty} \frac{(2l+1)}{e^{(\epsilon_l - \epsilon_0)/(k_B T)} - 1}$$

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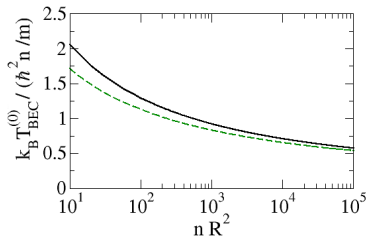
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$$\frac{k_B T_{\text{BEC}}}{\hbar^2 n / m} = \frac{-2\pi}{\ln \left[1 - e^{-\frac{\hbar^2 / (mR^2)}{k_B T_{\text{BEC}}} } \right]}$$



[AT, Salasnich, PRL **123**, 160403 (2019)]

Bose-Einstein condensation on the surface of a sphere

Uniform **interacting** bosons on the surface of the sphere

$$\Omega = -\frac{1}{\beta} \ln(\mathcal{Z}), \quad \mathcal{Z} = \int \mathcal{D}[\bar{\psi}, \psi] e^{-\frac{S[\bar{\psi}, \psi]}{\hbar}},$$

with $S[\bar{\psi}, \psi] = \int_0^{\beta\hbar} d\tau \int_0^{2\pi} d\varphi \int_0^\pi d\theta R^2 \sin\theta \mathcal{L}(\bar{\psi}, \psi)$, and

$$\mathcal{L} = \bar{\psi}(\theta, \varphi, \tau) \left(\hbar \partial_\tau + \frac{\hat{L}^2}{2mR^2} - \mu \right) \psi(\theta, \varphi, \tau) + \frac{g_0}{2} |\psi(\theta, \varphi, \tau)|^4.$$

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Bogoliubov theory:

$$\psi(\theta, \varphi, \tau) = \psi_0 + \eta(\theta, \varphi, \tau)$$

$$\eta(\theta, \varphi, \tau)$$



Bogoliubov spectrum

$$E_l^B = \sqrt{\epsilon_l(\epsilon_l + 2\mu)}$$

$$\psi_0$$



Equation of state of a spherical Bose gas

Implementing the Bogoliubov theory, we calculate

▷ Ω

▷ T_{BEC}

▷ n_0/n

▷ $n = \partial_\mu[\Omega/(4\pi R^2)]$

*: [AT, Pelster, Salasnich, arXiv:2104.04585], [AT, submitted to PRA]

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$$\triangleright n = \partial_\mu [\Omega / (4\pi R^2)]$$

In particular, we regularize the grand potential Ω through the analysis of scattering theory*. The number density reads:

$$n = \frac{m\mu}{4\pi\hbar^2} \ln \left\{ \frac{4\hbar^2 [1 - \alpha(\mu)]}{m\mu a_s^2 e^{2\gamma+1+\alpha(\mu)}} \right\} + \frac{1}{4\pi R^2} \sum_{l=1}^{\infty} \sum_{m_l=-l}^l \frac{\epsilon_l}{E_l^B} \frac{1}{e^{\beta E_l^B} - 1},$$

where we introduce the positive function $\alpha(\mu) = 1 - \frac{\mu}{\mu + E_1^B + \epsilon_1}$.

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Complete description of the quantum statistical properties of a weakly-interacting **spherical Bose gas**

*: [AT, Pelster, Salasnich, arXiv:2104.04585], [AT, submitted to PRA]

Outline

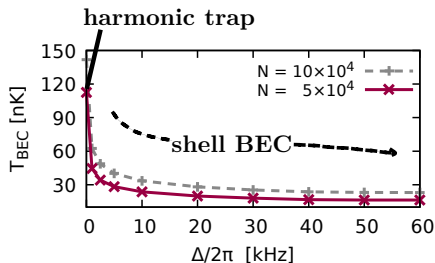
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Bose-Einstein condensation in ellipsoidal bubbles

In [AT, Cinti, Salasnich, PRL 125, 010402 (2020)], we modeled the microgravity experiments ([arXiv:2108.05880])

For the **realistic** trap parameters:

T_{BEC} drops quickly with $\Delta \propto$ shell area



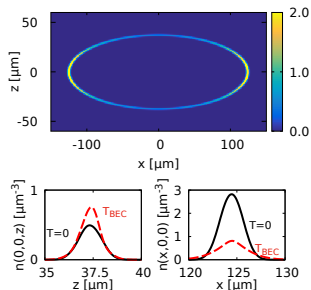
$N \sim 10^5$, $T_{BEC} \sim 30$ nK

Difficult to reach fully-condensate regime...

⇒ **Finite-temperature** properties are highly relevant

Density distribution

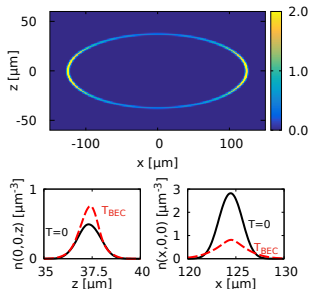
Condensate vs thermal density



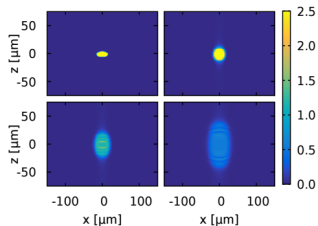
[AT, Cinti, Salasnich, PRL **125**, 010402 (2020)]

Density distribution and free expansion

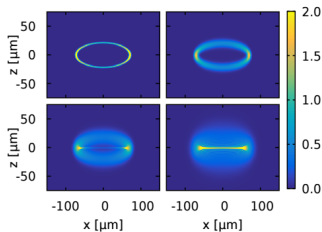
Condensate vs thermal density



Harmonic trap



Bubble trap



[AT, Cinti, Salasnich, PRL **125**, 010402 (2020)]

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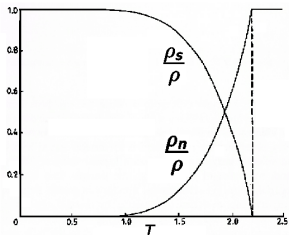
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Superfluidity: flow without friction

Landau two-fluid model:

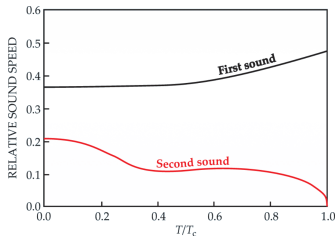


[Landau J. Phys. (USSR) 5, 71 (1941)]

predicts two **sound modes**

$$c_{1,2} = \left[\frac{v_A^2 + v_L^2}{2} \pm \sqrt{\left(\frac{v_A^2 + v_L^2}{2} \right)^2 - v_L^2 v_T^2} \right]^{1/2},$$

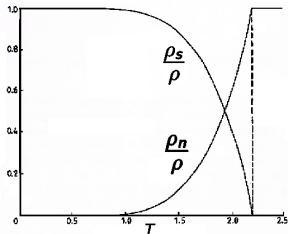
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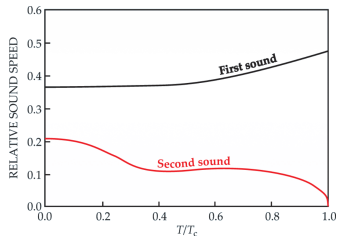


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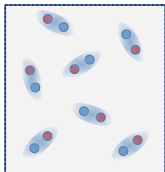
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superfluid transition in 2D?

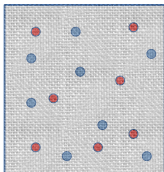
Berezinskii-Kosterlitz-Thouless transition (BKT)

$$T < T_{\text{BKT}}$$



superfluid

$$T > T_{\text{BKT}}$$

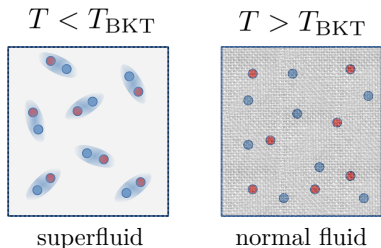


normal fluid

superfluid transition in 2D

[Nelson, Kosterlitz, PRL **39**, 1201
(1977)]

Berezinskii-Kosterlitz-Thouless transition (BKT)



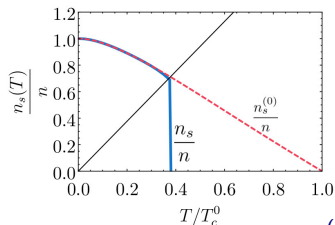
superfluid transition in 2D
[Nelson, Kosterlitz, PRL **39**, 1201
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RG equations of a flat superfluid

$$\frac{dK^{-1}(\ell)}{d\ell} = -4\pi^3 y^2(\ell)$$
$$\frac{dy(\ell)}{d\ell} = [2 - \pi K(\ell)] y(\ell)$$

Adimensional parameters

$$K(\ell) = \frac{\hbar^2 n_s(\ell)}{mk_B T}; \quad y(\ell) \sim e^{-\beta\mu_v(\ell)};$$



→ From bare $n_s(\ell=0) = n_s^{(0)}$
to renormalized $n_s = n_s(\ell=\infty)$

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BKT transition in spherical superfluids

RG equations of a spherical superfluid

$$\frac{dK^{-1}(\theta)}{d\ell(\theta)} = -4\pi^3 y^2(\theta)$$
$$\frac{dy(\theta)}{d\ell(\theta)} = [2 - \pi K(\theta)] y(\theta)$$

RG scale?

$$\ell(\theta) = \ln[2R \sin(\theta/2)/\xi]$$

Distance between vortices:

$$2R \sin(\theta/2) \in [\xi, 2R]$$

...in 3D space

BKT transition in spherical superfluids

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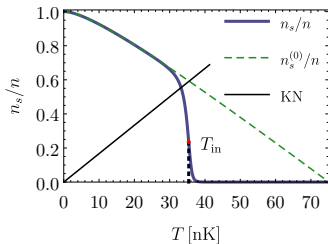
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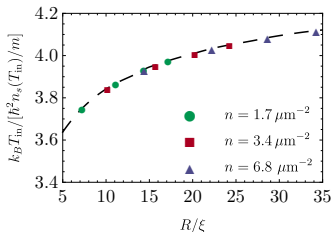
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...in 3D space

Finite system size \Rightarrow
smooth vanishing of n_s



for different n , shell widths
 T_{in} scales with R/ξ



BKT transition in spherical superfluids

RG equations of a spherical superfluid

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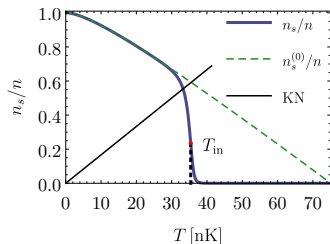
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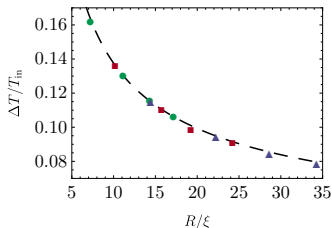
...in 3D space

Finite system size \Rightarrow
smooth vanishing of n_s



for different n , shell widths

$$\Delta T/T_{in} \propto \ln^{-2}(R/\xi)$$



[AT, Pelster, Salasnich, arXiv:2104.04585]

Hydrodynamic modes in spherical superfluids

Given the **equation of state** **and** the **superfluid density**, we extend the Landau two-fluid model to the spherical case.

Hydrodynamic modes in spherical superfluids

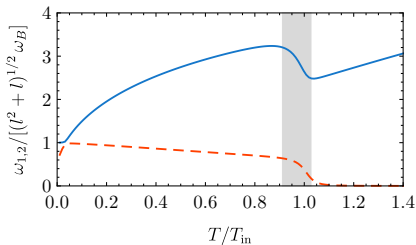
Given the **equation of state** **and** the **superfluid density**, we extend the Landau two-fluid model to the spherical case.

Frequencies of the hydrodynamic modes:

$$\omega_{1,2}^2 = \left[\frac{l(l+1)}{R^2} \right] \left[\frac{v_A^2 + v_L^2}{2} \pm \sqrt{\left(\frac{v_A^2 + v_L^2}{2} \right)^2 - v_L^2 v_T^2} \right]$$

ω_1, ω_2 are the **main quantitative probe of BKT physics**

$$v_{\{A,T\}} = \sqrt{\left(\frac{\partial P}{\partial \rho} \right)_{\{\xi,T\}}}, \quad v_L = \sqrt{\frac{\rho_s T \xi^2}{\rho_n \tilde{c}_V}}$$



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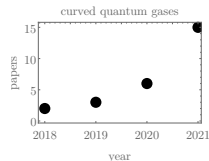
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- Curvature in quantum gases (and in cond. mat.): a new research direction.

The scientific community has just started exploring shell-shaped BECs

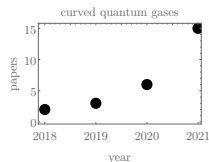


Conclusions

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The scientific community has just started exploring shell-shaped BECs

- in spherical condensates: curvature \approx finite-size, but consequences on superfluidity
- interesting perspectives with ellipsoidal shells



Thank you for your attention!

References



A. Tononi, F. Cinti, and L. Salasnich, *Quantum Bubbles in Microgravity*, Physical Review Letters **125**, 010402 (2020).



A. Tononi and L. Salasnich, *Bose-Einstein Condensation on the Surface of a Sphere*, Physical Review Letters **123**, 160403 (2019).



A. Tononi, A. Pelster, and L. Salasnich, arXiv:2104.04585