

# Quantum Bubbles in Microgravity

Andrea Tononi

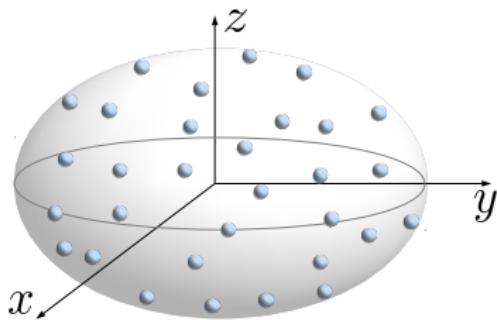
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This presentation on [www.andreatononi.com](http://www.andreatononi.com)

# Quantum bubbles

Bose-Einstein condensate on a thin ellipsoidal shell

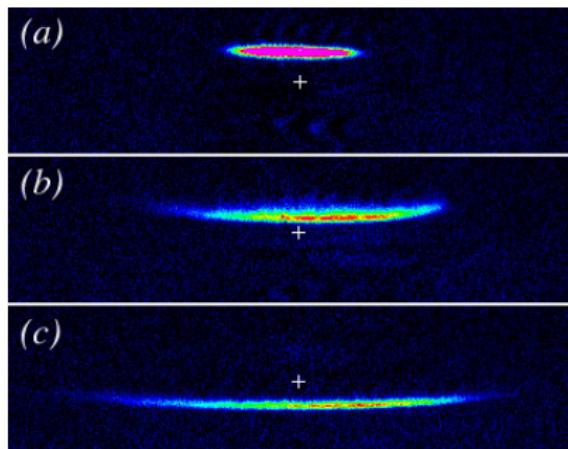


For strong radial confinement:  
**2D curved superfluid**

- finite-size 2D condensate
- topology (and BKT)
- quantized vortices
- free expansion

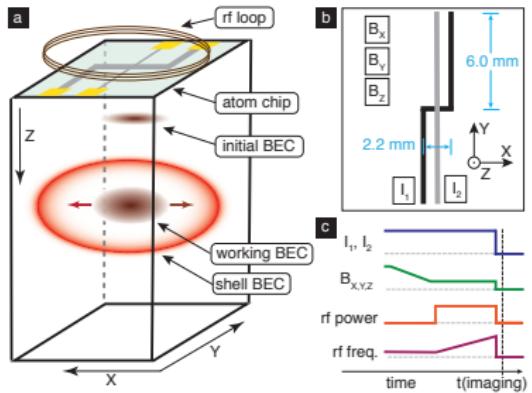
# Why microgravity?

Bubble-trap on Earth...



[Colombe, *et al.*, EPL 2004]

NASA-JPL Cold Atom Laboratory

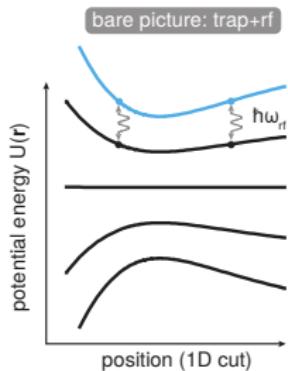


[Lundblad, *et al.*, npj Microgravity 2019]

# Outline

- ▷ Introduction on bubble trapping
- ▷ Bose-Einstein condensation on the surface of a sphere
- ▷ Shell-shaped condensates: statics and dynamics
- ▷ Summary and outlook

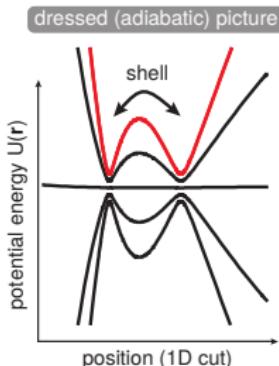
# Bubble-trapping: experimental realization



Alkali-metal atoms with total angular momentum  $F = 2$ .

- + Magnetic field  $\mathbf{B}(\vec{r}) \implies$  space-dependent Zeeman splitting with  $m_F = \{\pm 2, \pm 1, 0\} \implies$  space-dependent bare potentials  $u(\vec{r})$

- + Radiofrequency field  $\mathbf{B}_{\text{rf}}(\vec{r}, t) \implies$  bubble-trap in the dressed picture (old  $m_F$  bad quantum number)



[Lundblad, et al., npj Microgravity 2019]

## Bubble-trap

$$U(\vec{r}) = M_F \sqrt{\left[ \sum_i m \omega_i^2 x_i^2 / 2 - \hbar \Delta \right]^2 + (\hbar \Omega)^2},$$

$\omega_i$ : frequencies of the bare harmonic trap

$\Delta$ : detuning from the resonant frequency

$\Omega$ : Rabi frequency between coupled levels

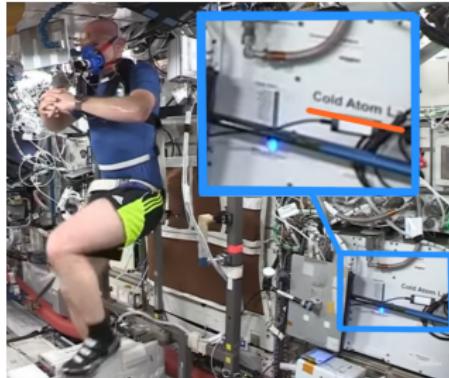
Minimum for  $\omega_x^2 x^2 + \omega_y^2 y^2 + \omega_z^2 z^2 = 2\hbar\Delta/m$ .

[Zobay, Garraway, PRL 2001]

If gravity is included the **atoms will pool on the bottom of the trap!**

$$U(\vec{r}) = M_F \sqrt{\left[ \sum_i m \omega_i^2 x_i^2 / 2 - \hbar \Delta \right]^2 + (\hbar \Omega)^2} + \underline{mgz},$$

→ Experiments on the NASA-JPL **Cold Atom Laboratory** on the International Space Station (PI Nathan Lundblad).



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# Bose-Einstein condensation on the surface of a sphere

Noninteracting bosons.

Quantized energy  $\varepsilon_I = \frac{\hbar^2}{2mR^2} I(I+1)$ , with degeneracy  $2I+1$ .

In the Bose-condensed phase, we can set  $\mu = 0$  and

$$N = N_0 + \sum_{I=1}^{+\infty} \frac{2I+1}{e^{\varepsilon_I/(k_B T)} - 1}$$

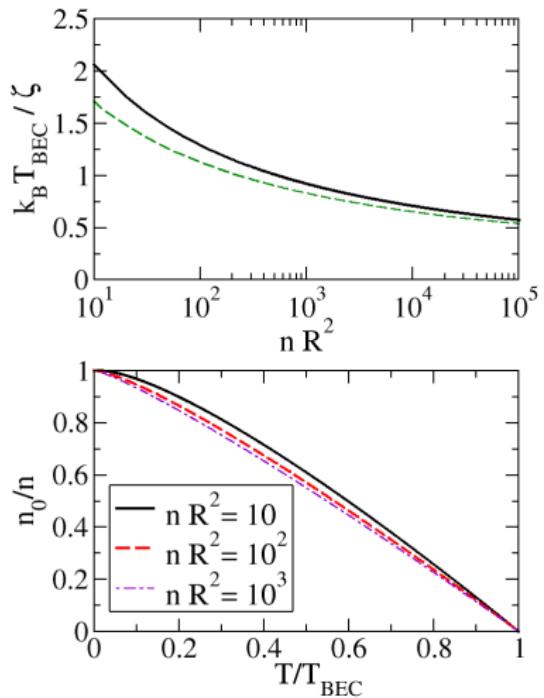
when  $N_0 = 0 \implies T = T_{\text{BEC}}$

# BEC on a sphere: noninteracting case

$$k_B T_{\text{BEC}} =$$

$$\frac{\frac{2\pi\hbar^2}{m}n}{\frac{\beta_{\text{BEC}}\hbar^2}{mR^2} - \ln(e^{\beta_{\text{BEC}}\hbar^2/mR^2} - 1)}$$

$$\frac{n_0}{n} = 1 - \frac{1 - \frac{mR^2}{\hbar^2\beta} \ln(e^{\beta\hbar^2/mR^2} - 1)}{1 - \frac{mR^2}{\hbar^2\beta_{\text{BEC}}} \ln(e^{\beta_{\text{BEC}}\hbar^2/mR^2} - 1)}$$



[AT, Salasnich, PRL 123, 160403 (2019)]

## BEC on a sphere: interacting case

The grand potential  $\Omega = -\beta^{-1} \ln(\mathcal{Z})$ , where  $\mathcal{Z}$  is the grand canonical partition function

$$\mathcal{Z} = \int \mathcal{D}[\bar{\psi}, \psi] e^{-S[\bar{\psi}, \psi]/\hbar},$$

where

$$S[\bar{\psi}, \psi] = \int_0^{\beta\hbar} d\tau \int_0^{2\pi} d\varphi \int_0^\pi \sin(\theta) d\theta R^2 \mathcal{L}(\bar{\psi}, \psi)$$

is the Euclidean action, and

$$\mathcal{L} = \bar{\psi}(\theta, \varphi, \tau) \left( \hbar \partial_\tau + \frac{\hbar^2}{2mR^2} - \mu \right) \psi(\theta, \varphi, \tau) + \frac{g}{2} |\psi(\theta, \varphi, \tau)|^4$$

is the Euclidean Lagrangian.

## BEC on a sphere: interacting case

In the Bose-condensed phase

$$\psi(\theta, \varphi, \tau) = \psi_0 + \eta(\theta, \varphi, \tau)$$

Expanding up to quadratic order, expanding with spherical harmonics, and performing functional integration we get

$$\begin{aligned}\Omega(\mu, \psi_0^2) &= 4\pi R^2 \left( -\mu\psi_0^2 + g\psi_0^4/2 \right) + \frac{\alpha}{2} \sum_{l=1}^{\infty} \sum_{m_l=-l}^l E_l(\mu, \psi_0^2) \\ &\quad + \frac{\alpha}{\beta} \sum_{l=1}^{\infty} \sum_{m_l=-l}^l \ln \left( 1 - e^{-\beta E_l(\mu, \psi_0^2)} \right) + o(\alpha^2),\end{aligned}$$

with  $E_l(\mu, \psi_0^2) = \sqrt{(\epsilon_l - \mu + 2g\psi_0^2)^2 - g^2\psi_0^4}$ .

## BEC on a sphere: interacting case

Applying Variational Perturbation Theory we can calculate the critical temperature of the interacting system

$$k_B T_{\text{BEC}} = \frac{\frac{2\pi\hbar^2 n}{m} - \frac{gn}{2}}{\frac{\hbar^2 \beta_{\text{BEC}}}{2mR^2} \left( 1 + \sqrt{1 + \frac{2gmnR^2}{\hbar^2}} \right) - \ln \left( e^{\frac{\hbar^2 \beta_{\text{BEC}}}{mR^2}} \sqrt{1 + \frac{2gmnR^2}{\hbar^2}} - 1 \right)}.$$

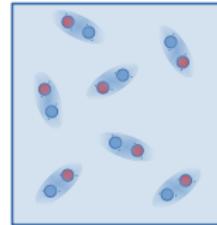
and the condensate fraction

$$\begin{aligned} \frac{n_0}{n} &= 1 - \frac{mg}{4\pi\hbar^2} - \frac{1}{4\pi R^2 n} \left[ 1 + \sqrt{1 + \frac{2gmnR^2}{\hbar^2}} \right] \\ &+ \frac{mk_B T}{2\pi\hbar^2 n} \ln \left( e^{\frac{\hbar^2}{mR^2 k_B T}} \sqrt{1 + (2gmnR^2/\hbar^2)} - 1 \right). \end{aligned}$$

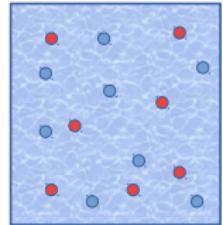
$R \rightarrow \infty$ :  $T_{\text{BEC}} \rightarrow 0$ , Schick result for quantum depletion.

The unbinding of vortex-antivortex dipoles at  $T = T_{\text{BKT}}$  destroys the quasi long-range order.

$$T < T_{\text{BKT}}$$



$$T > T_{\text{BKT}}$$



Kosterlitz-Nelson criterion on the sphere [Ovrut, Thomas PRD 1991]

$$k_B T_{\text{BKT}} = \frac{\pi}{2} \frac{\hbar^2 n_s(T_{\text{BKT}})}{m}$$

with the superfluid density  $n_s(T)$  as

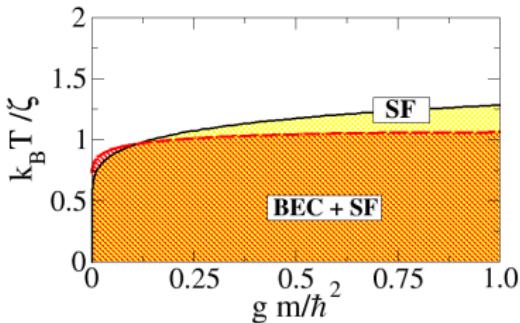
$$n_s = n - \frac{1}{k_B T} \int_1^{+\infty} \frac{dl (2l+1)}{4\pi R^2} \frac{\hbar^2(l^2+l)}{2mR^2} \frac{e^{E_l^B/(k_B T)}}{(e^{E_l^B/(k_B T)} - 1)^2}.$$

# BEC on a sphere: interacting case

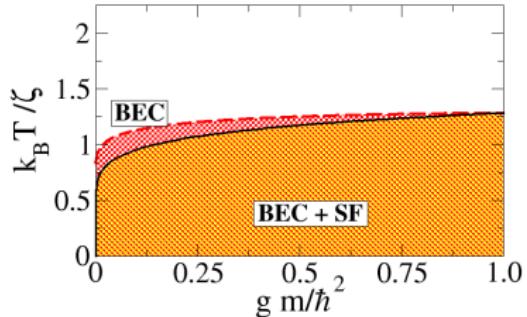
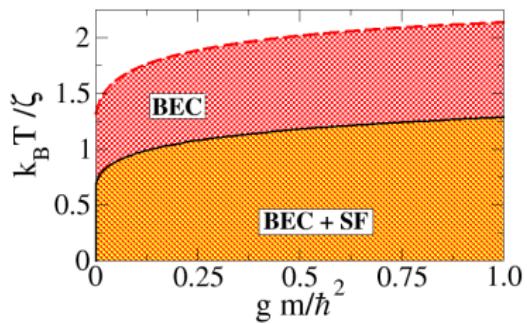
BEC transition (red dashed)  
BKT=SF transition (black)

Plots:  $nR^2 =$        $10^2$   
                 $10^5 \quad 10^4$

Usual 2D picture (thermodyn. limit)



Region of BEC only



[AT, Salasnich, PRL 123, 160403 (2019)]

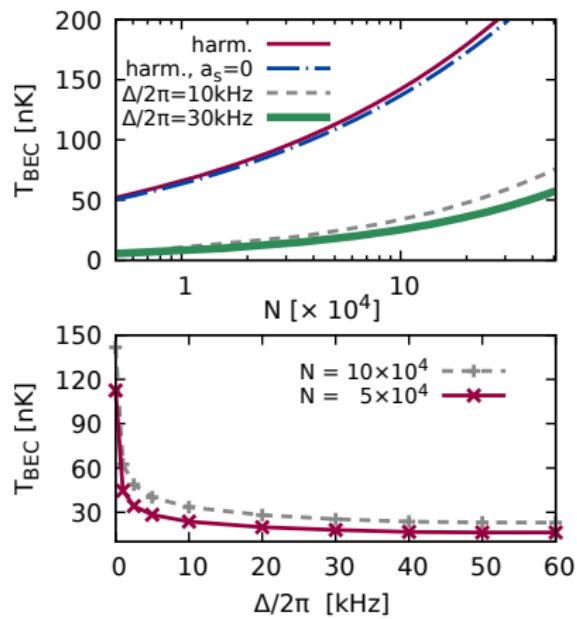
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# Shell-shaped (ellipsoidal) condensates

For the **realistic** trap parameters of  
NASA-JPL CAL experiment:

$$T_{BEC}^{bubble\ trap} \ll T_{BEC}^{harmonic\ trap}$$

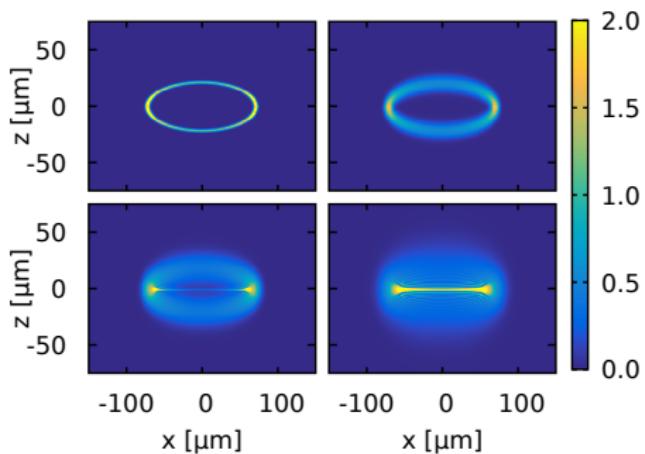


[AT, Cinti, Salasnich, arxiv:1912.07297]

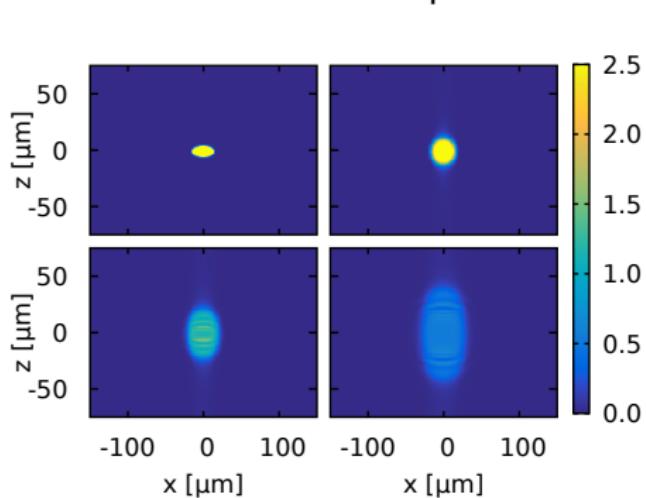
(using Hartree-Fock theory [Giorgini, et al., J. Low T. Phys., **109**, 309 (1997)] )

# Free expansion

Bubble trap

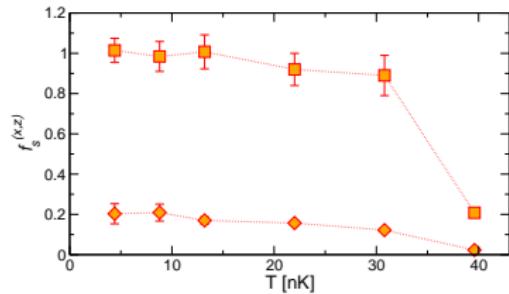
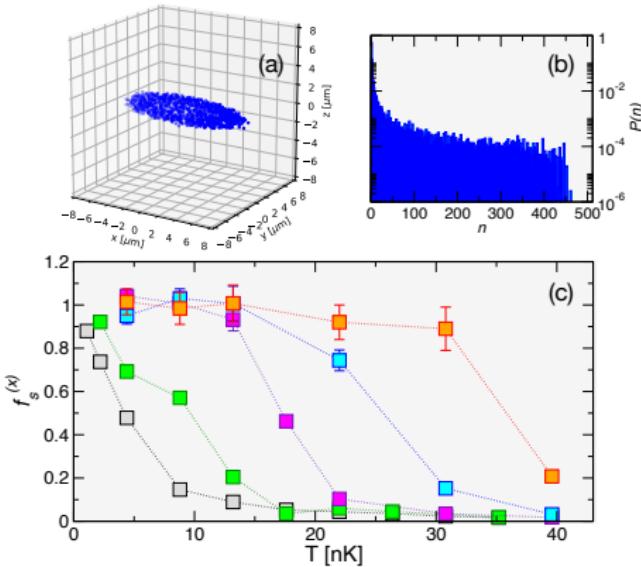


Harmonic trap



[AT, Cinti, Salasnich, arxiv:1912.07297]

# Path Integral Monte Carlo simulations



anisotropic  $f_s^{(x,z)}$  from  
nonclassical moment of inertia

[AT, Cinti, Salasnich, arxiv:1912.07297]

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# Summary and outlook

## SPHERICAL TRAP

- ◊ Calculation of  $T_{\text{BEC}}$ ,  $T_{\text{BKT}}$ , and  $n_0/n$  for spherical trap

## ELLIPSOIDAL TRAP

- ◊  $T_{\text{BEC}}$  and realistic experimental description
- ◊ Free expansion
- ◊ Superfluidity studied with Quantum Monte Carlo

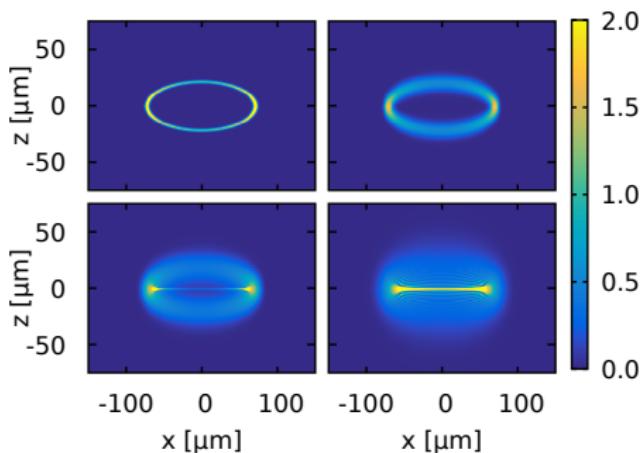
## Future projects

Vortices on a compact superfluid

Extension of the analytical calculations

# Quantum bubbles in microgravity

Andrea Tononi



# References

-  A. Tononi, L. Salasnich, *Bose-Einstein Condensation on the Surface of a Sphere*, Phys. Rev. Lett. **123**, 160403 (2019).
-  A. Tononi, F. Cinti, L. Salasnich, *Quantum bubbles in microgravity*, arxiv:1912.07297
-  N. Lundblad, R. A. Carollo, C. Lannert, et al. *Shell potentials for microgravity Bose-Einstein condensates*, npj Microgravity **5**, 30 (2019).