# Quantum Bubbles in Microgravity



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## Quantum bubbles

#### Bose-Einstein condensate on a thin ellipsoidal shell



For strong radial confinement: **2D curved superfluid** 

finite-size 2D condensatetopology (and BKT)

- quantized vortices
- o free expansion

## Why microgravity?

Bubble-trap on Earth...



#### NASA-JPL Cold Atom Laboratory



#### [Colombe, et al., EPL 2004]

[Lundblad, et al., npj Microgravity 2019]



- Introduction on bubble trapping
- ▷ Bose-Einstein condensation on the surface of a sphere
- Shell-shaped condensates: statics and dynamics
- Summary and outlook

## Bubble-trapping: experimental realization



Alkali-metal atoms with total angular momentum F = 2.

+ Magnetic field  $\mathbf{B}(\vec{r}) \implies$ space-dependent Zeeman splitting with  $m_F = \{\pm 2, \pm 1, 0\} \implies$ space-dependent bare potentials  $u(\vec{r})$ 

+ Radiofrequency field  $\mathbf{B}_{rf}(\vec{r}, t) \implies$ bubble-trap in the dressed picture (old  $m_F$  bad quantum number)

[Lundblad, et al., npj Microgravity 2019]



#### Bubble-trap

$$U(\vec{r}) = M_F \sqrt{\left[\sum_{i} m \omega_i^2 x_i^2 / 2 - \hbar \Delta\right]^2 + (\hbar \Omega)^2},$$

- $\omega_i:$  frequencies of the bare harmonic trap
- $\Delta$ : detuning from the resonant frequency
- $\Omega$ : Rabi frequency between coupled levels

Minimum for 
$$\omega_x^2 x^2 + \omega_y^2 y^2 + \omega_z^2 z^2 = 2\hbar\Delta/m$$
.

[Zobay, Garraway, PRL 2001]

If gravity is included the atoms will pool on the bottom of the trap!

$$U(\vec{r}) = M_F \sqrt{\left[\sum_{i} m \omega_i^2 x_i^2 / 2 - \hbar \Delta\right]^2 + (\hbar \Omega)^2 + \underline{mgz}},$$

→ Experiments on the Cold Atom Lab on the International Space Station (PI Nathan Lundblad). [Aveline, *et al.*, Nature **582**, 193 (2020)]







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#### Bose-Einstein condensation on the surface of a sphere

Noninteracting bosons.

Quantized energy  $\varepsilon_I = \frac{\hbar^2}{2mR^2} I(I+1)$ , with degeneracy 2I + 1.

In the Bose-condensed phase, we can set  $\mu = 0$  and

$$N = N_0 + \sum_{l=1}^{+\infty} \frac{2l+1}{e^{\varepsilon_l/(k_B T)} - 1}$$

when  $N_0 = 0 \implies T = T_{BEC}$ 

#### BEC on a sphere: **noninteracting case**



## BEC on a sphere: interacting case

The grand potential  $\Omega = -\beta^{-1} \ln(\mathcal{Z})$ , where  $\mathcal{Z}$  is the grand canonical partition function

$$\mathcal{Z} = \int \mathcal{D}[\bar{\psi}, \psi] \; e^{-\mathcal{S}[\bar{\psi}, \psi]/\hbar},$$

where

$$S[\bar{\psi},\psi] = \int_0^{\beta\hbar} d\tau \, \int_0^{2\pi} d\varphi \, \int_0^{\pi} \sin(\theta) \, d\theta \, R^2 \, \mathcal{L}(\bar{\psi},\psi)$$

is the Euclidean action, and

$$\mathcal{L} = \bar{\psi}(\theta, \varphi, \tau) \bigg( \hbar \partial_{\tau} + \frac{L^2}{2mR^2} - \mu \bigg) \psi(\theta, \varphi, \tau) + \frac{g}{2} |\psi(\theta, \varphi, \tau)|^4$$

is the Euclidean Lagrangian.

In the Bose-condensed phase

$$\psi(\theta,\varphi,\tau) = \psi_0 + \eta(\theta,\varphi,\tau)$$

Expanding up to quadratic order, expanding with spherical harmonics, and performing functional integration we get

$$\begin{split} \Omega(\mu,\psi_0^2) &= 4\pi R^2 \big(-\mu\psi_0^2 + g\psi_0^4/2\big) + \frac{\alpha}{2} \sum_{l=1}^{\infty} \sum_{m_l=-l}^{l} E_l(\mu,\psi_0^2) \\ &+ \frac{\alpha}{\beta} \sum_{l=1}^{\infty} \sum_{m_l=-l}^{l} \ln\left(1 - e^{-\beta E_l(\mu,\psi_0^2)}\right) + o(\alpha^2), \end{split}$$

with 
$$E_I(\mu, \psi_0^2) = \sqrt{(\epsilon_I - \mu + 2g\psi_0^2)^2 - g^2\psi_0^4}$$
.

## BEC on a sphere: interacting case

Applying Variational Perturbation Theory we can calculate the critical temperature of the interacting system

$$k_{B}T_{BEC} = \frac{\frac{2\pi\hbar^{2}n}{m} - \frac{gn}{2}}{\frac{\hbar^{2}\beta_{BEC}}{2mR^{2}}\left(1 + \sqrt{1 + \frac{2gmnR^{2}}{\hbar^{2}}}\right) - \ln\left(e^{\frac{\hbar^{2}\beta_{BEC}}{mR^{2}}\sqrt{1 + \frac{2gmnR^{2}}{\hbar^{2}}} - 1\right)}.$$

and the condensate fraction

$$\frac{n_0}{n} = 1 - \frac{mg}{4\pi\hbar^2} - \frac{1}{4\pi R^2 n} \left[ 1 + \sqrt{1 + \frac{2gmnR^2}{\hbar^2}} \right] + \frac{mk_B T}{2\pi\hbar^2 n} \ln\left(e^{\frac{\hbar^2}{mR^2 k_B T}\sqrt{1 + (2gmnR^2/\hbar^2)}} - 1\right).$$

 $R \rightarrow \infty:~ T_{\mathsf{BEC}} \rightarrow 0,$  Schick result for quantum depletion.

The unbinding of vortex-antivortex dipoles at  $T = T_{BKT}$  destroys the quasi long-range order.

$$T < T_{\rm BKT} \qquad T > T_{\rm BKT}$$

Kosterlitz-Nelson criterion on the sphere [Ovrut, Thomas PRD 1991]

$$k_B T_{BKT} = \frac{\pi}{2} \frac{\hbar^2 n_s(T_{BKT})}{m}$$

with the superfluid density  $n_s(T)$  as

$$n_{s} = n - \frac{1}{k_{B}T} \int_{1}^{+\infty} \frac{dl(2l+1)}{4\pi R^{2}} \frac{\hbar^{2}(l^{2}+l)}{2mR^{2}} \frac{e^{E_{l}^{B}/(k_{B}T)}}{(e^{E_{l}^{B}/(k_{B}T)}-1)^{2}}.$$

## BEC on a sphere: interacting case

BEC transition (red dashed) BKT=SF transition (black)

Plots:  $nR^2 = 10^2$  $10^5 \ 10^4$ 

Usual 2D picture (thermodyn. limit)



Region of BEC only



[AT, Salasnich, PRL 123, 160403 (2019)]



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## Shell-shaped (ellipsoidal) condensates

# For the realistic trap parameters of NASA-JPL CAL experiment:

 $T_{BEC}^{bubble\ trap} \ll T_{BEC}^{harmonic\ trap}$ 



[AT, Cinti, Salasnich, arxiv:1912.07297]

(using Hartree-Fock theory [Giorgini, et al., J. Low T. Phys., 109, 309 (1997)] )

#### Free expansion

Bubble trap

Harmonic trap



[AT, Cinti, Salasnich, arxiv:1912.07297]

#### Path Integral Monte Carlo simulations



anisotropic  $f_s^{(x,z)}$  from nonclassical moment of inertia

[AT, Cinti, Salasnich, arxiv:1912.07297]



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Summary and outlook

#### SPHERICAL TRAP

 $\diamond\,$  Calculation of  $T_{\rm BEC},\,T_{\rm BKT},$  and  $\mathit{n_0/n}$  for spherical trap

ELLIPSOIDAL TRAP

- $\diamond$  T<sub>BEC</sub> and realistic experimental description
- ◊ Free expansion
- Superfluidity studied with Quantum Monte Carlo

#### Future projects

Vortices on a compact superfluid

Extension of the analytical calculations

# Quantum bubbles in microgravity

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#### References

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