Equilibrium and nonequilibrium physics of bubble-trapped condensates

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In collaboration with F. Cinti, A. Pelster, L. Salasnich. This presentation is on

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A mini-talk/poster on

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Bubble-trapped BECs



curved and hollow 2D Bose gas A mini-talk/poster on

Bubble-trapped BECs



curved and hollow 2D Bose gas Contents:

- Experiments overview and numerical results
- $\circ \text{ Spherical shell} \longrightarrow T_{BEC}, \\ \text{equation of state}$
- BKT transition
- Hydrodynamic excitations

Experimentally realizable Bubble-trap...



[Lundblad et al., npj Microgravity 5, 30 (2019)]

...on Earth



[Colombe et al., EPL 67, 593 (2004)]

Experimentally realizable...in microgravity Bubble-trap... ...on Earth





[Colombe et al., EPL 67, 593 (2004)]

⇒ Experiments on NASA Cold Atom Lab

[Carollo *et al.*, arXiv:2108.05880] [Aveline *et al.*, Nature **582**, 193 (2020)]



Numerical results — inflating the ellipsoidal bubble

Bubble-trap: $U(\vec{r}) = M_F \sqrt{\left[\sum_i m \omega_i^2 x_i^2 / 2 - \hbar \Delta\right]^2 + (\hbar \Omega)^2},$ [Zobay, Garraway, Phys. Rev. Lett. **86**, 1195 (2001)]

For the realistic trap parameters of NASA-JPL CAL experiment:

 T_{BEC} drops quickly with $\Delta \propto$ shell area



[AT, Cinti, Salasnich, PRL 125, 010402 (2020)]

Difficult to reach fully-condensate regime...

 \Rightarrow Finite-temperature properties and BKT physics are highly relevant

Numerical results — density at T = 0 and at T_{BEC}

Number density as a probe of the system temperature

Gross-Pitaevskii equation at T = 0:

$$\left[-rac{\hbar^2
abla^2}{2m}+U(ec{r})+g|\psi|^2
ight]\psi=\mu\psi$$

Hartree-Fock at T_{BEC} :

$$n(\vec{r}) = \int \frac{d^3p}{e^{(E^{HF}(\vec{p},\vec{r}))/(k_B T_{BEC})} - 1}$$
$$E^{HF}(\vec{p},\vec{r}) = p^2/(2m) + U(\vec{r}) - \mu + 2gn(\vec{r})$$



[AT, Cinti, Salasnich, PRL 125, 010402 (2020)]

Numerical results — free expansion



[AT, Cinti, Salasnich, PRL 125, 010402 (2020)]

Can we have analytical results?

Bose-Einstein condensation on the surface of a sphere

Can we have analytical results? \rightarrow spherical symmetry

Noninteracting case, single particle on a sphere of radius *R*: $\frac{\hat{L}^2}{2mR^2} \psi_{l,m_l}(\theta,\varphi) = \epsilon_l \psi_{l,m_l}(\theta,\varphi),$ with $\epsilon_l = \frac{\hbar^2}{2mR^2} l(l+1)$ and $m_l = -l, \dots, +l$.

Particle number at temperature T:

$$N = N_0 + \sum_{l=1}^{+\infty} \sum_{m_l=-l}^{+l} \frac{1}{e^{(\epsilon_l - \epsilon_0)/(k_B T)} - 1}$$

when $N_0 = 0 \implies T = T_{\text{BEC}}^{(0)}$

[AT, Salasnich, PRL 123, 160403 (2019)]

BEC on a sphere: **noninteracting case**

$$k_B T_{BEC}^{(0)} = \frac{\frac{2\pi\hbar^2}{m}n}{\frac{\beta_{BEC}^{(0)}\hbar^2}{mR^2} - \ln(e^{\beta_{BEC}^{(0)}\hbar^2/mR^2} - 1)}$$

$$\begin{aligned} \frac{n_0}{n} &= 1 - \\ \frac{1 - \frac{mR^2}{\hbar^2\beta} \ln\left(e^{\beta\hbar^2/mR^2} - 1\right)}{1 - \frac{mR^2}{\hbar^2\beta_{\text{BEC}}^{(0)}} \ln\left(e^{\beta_{\text{BEC}}^{(0)}\hbar^2/mR^2} - 1\right)} \end{aligned}$$



[AT, Salasnich, PRL 123, 160403 (2019)]

BEC on a sphere: interacting case

With Gaussian functional integration: interacting critical temperature

$$k_B T_{\text{BEC}} = \frac{\frac{2\pi\hbar^2 n}{m} - \frac{gn}{2}}{\frac{\hbar^2 \beta_{BEC}}{2mR^2} \left(1 + \sqrt{1 + \frac{2gmnR^2}{\hbar^2}}\right) - \ln\left(e^{\frac{\hbar^2 \beta_{BEC}}{mR^2}\sqrt{1 + \frac{2gmnR^2}{\hbar^2}}} - 1\right)}.$$

and the condensate fraction

$$\frac{n_0}{n} = 1 - \frac{mg}{4\pi\hbar^2} - \frac{1}{4\pi R^2 n} \left[1 + \sqrt{1 + \frac{2gmnR^2}{\hbar^2}} \right] \\ + \frac{mk_B T}{2\pi\hbar^2 n} \ln\left(e^{\frac{\hbar^2}{mR^2k_B T}\sqrt{1 + (2gmnR^2/\hbar^2)}} - 1\right).$$

 $R \rightarrow \infty$: $T_{\text{BEC}} \rightarrow 0$, flat-case quantum depletion [Schick, PRA 3, 1067 (1971)] [AT, Salasnich, PRL **123**, 160403 (2019)]

Equation of state

Recent calculation of the finite-temperature grand potential

$$\begin{split} &\frac{\Omega}{4\pi R^2} = -\frac{m\mu^2}{8\pi\hbar^2} \bigg[\ln \bigg(\frac{4\hbar^2}{m(E_1^B + \epsilon_1 + \mu)a^2 e^{2\gamma + 1}} \bigg) + \frac{1}{2} \bigg] \\ &+ \frac{mE_1^B}{8\pi\hbar^2} (E_1^B - \epsilon_1 - \mu) + \frac{1}{4\pi R^2} \frac{1}{\beta} \sum_{l=1}^{\infty} \sum_{m_l=-l}^{l} \ln \Big(1 - e^{-\beta E_l^B} \Big), \end{split}$$

[AT, Pelster, Salasnich, arXiv:2104.04585] [AT, under review]

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from which we calculate all the thermodynamic functions,

which will become useful to calculate hydrodynamic excitations...



[AT, Pelster, Salasnich, arXiv:2104.04585] [AT, under review]

BKT transition of a 2D spherical superfluid

Kosterlitz-Nelson analysis to describe the superfluid transition yields

RG equations as in flat case

$$egin{aligned} rac{d\mathcal{K}^{-1}(heta)}{d\ell(heta)} &= -4\pi^3 y^2(heta) \ rac{dy(heta)}{d\ell(heta)} &= \left[2 - \pi \mathcal{K}(heta)
ight] y(heta) \end{aligned}$$

RG scale? $\ell(\theta) = \ln[2R\sin(\theta/2)/\xi]$

Distance between vortices: $2R\sin(\theta/2) \in [\xi, 2R]...$

...but in 3D space!!

[AT, Pelster, Salasnich, arXiv:2104.04585]

BKT transition of a 2D spherical superfluid



Probing BKT physics with hydrodynamic modes

Response of a finite-temperature superfluid to a small perturbation:

Flat case: ordinary first and second sound (basis: plane waves $e^{i(kx-\omega t)}$) Shell BECs: hydrodynamic modes ω_1 , ω_2 (basis: spherical harmonics $\mathcal{Y}_l^{m_l} e^{i\omega t}$)



[Christodoulou, et al. Nature 594, 191 (2021)]



[AT, Pelster, Salasnich, arXiv:2104.04585]

The surface modes ω_1 , ω_2 are the main quantitative probe of BKT physics

In conclusion

We investigate equilibrium and nonequilibrium properties of **2D bubble-trapped quantum gases**, either ellipsoidal, or spherical.

Interesting topics:

- BKT transition
- vortices

Understanding, probing, controlling the relation quantum many-body physics ↔ curvature

Thank you for your attention!

References



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A. Tononi, A. Pelster, and L. Salasnich, arXiv:2104.04585