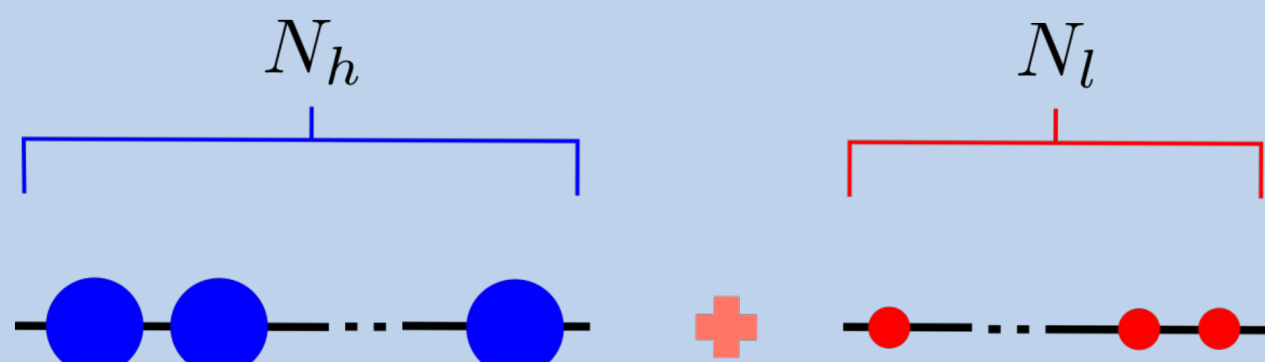


Self-binding of one-dimensional fermionic mixtures

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System

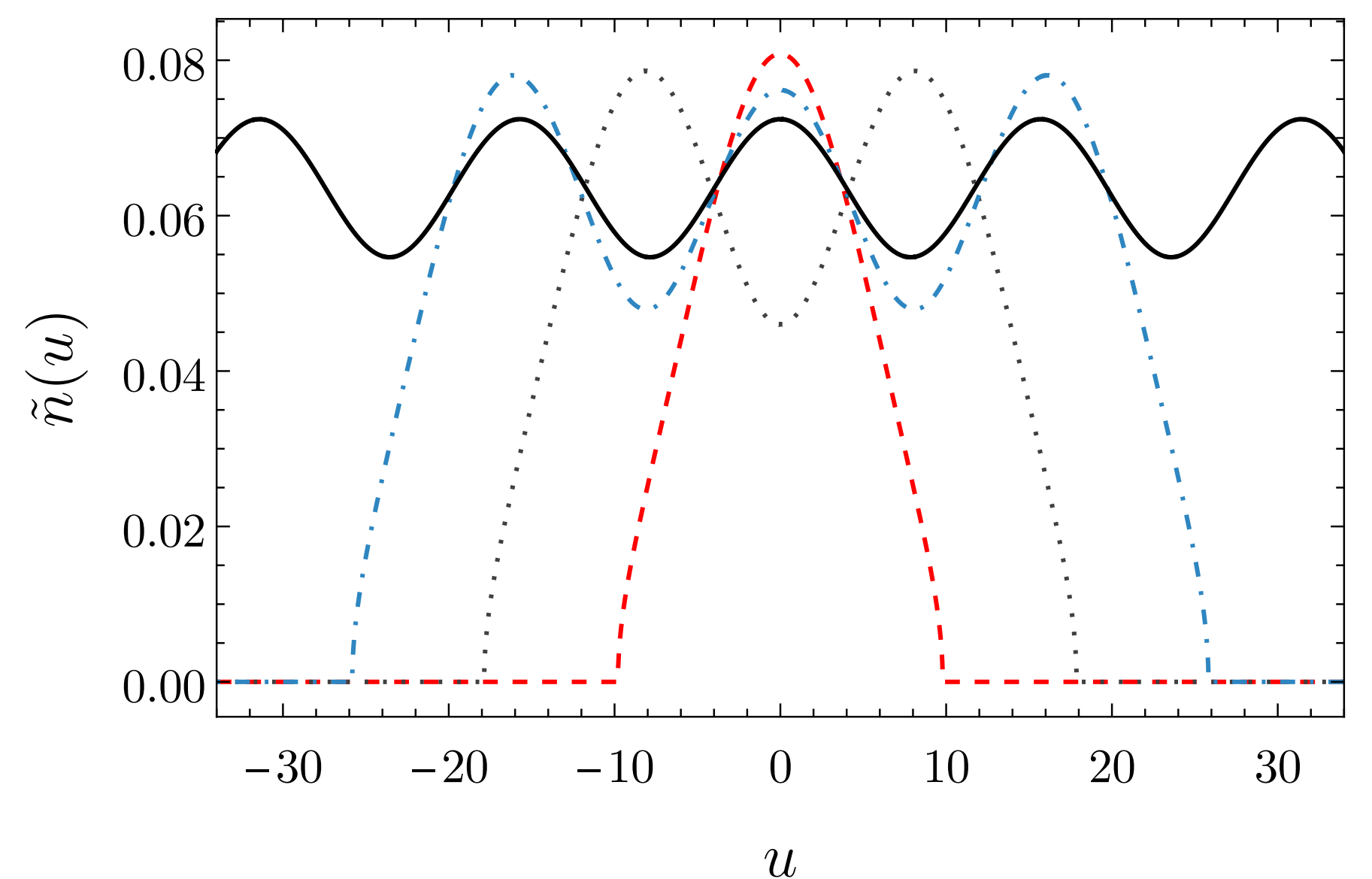
- N_h heavy fermions of mass M
- N_l light fermions of mass m
- Inter-species attraction ($g < 0$)
- No intra-species interaction



Competition between the Pauli repulsion and the interspecies attraction

$$\hat{H} = \int \left(-\frac{\hat{\Psi}^\dagger \partial_x^2 \hat{\Psi}}{2M} - \frac{\hat{\phi}^\dagger \partial_x^2 \hat{\phi}}{2m} + g \hat{\Psi}^\dagger \hat{\phi}^\dagger \hat{\Psi} \hat{\phi} \right) dx$$

Question: what is the ground state of this system?

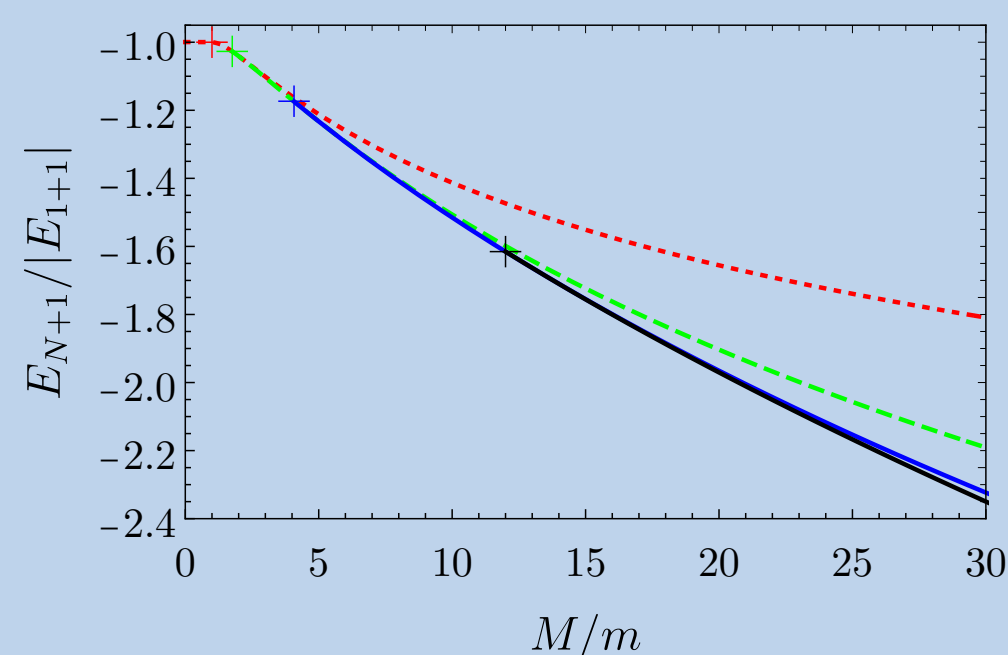


N+1 [1]

First off, let us start by taking $N_l = 1$ and $N_h = N$

Finite N

We solve exactly the **Skorniakov-Ter Martirosian** equation and find the bound state energies for up to $N = 5$



Critical mass-ratio for binding:

$$\begin{aligned} (M/m)_{2+1} &= 1 \\ (M/m)_{3+1} &= 1.76 \\ (M/m)_{4+1} &= 4.2 \\ (M/m)_{5+1} &= 12.0 \pm 0.5 \end{aligned}$$

Large N limit

We use a mean-field method, namely the **Thomas-Fermi** approximation for the kinetic energy of the heavy atom (in the limit where $N \gg 1$). The idea is to then write and minimize the grand potential

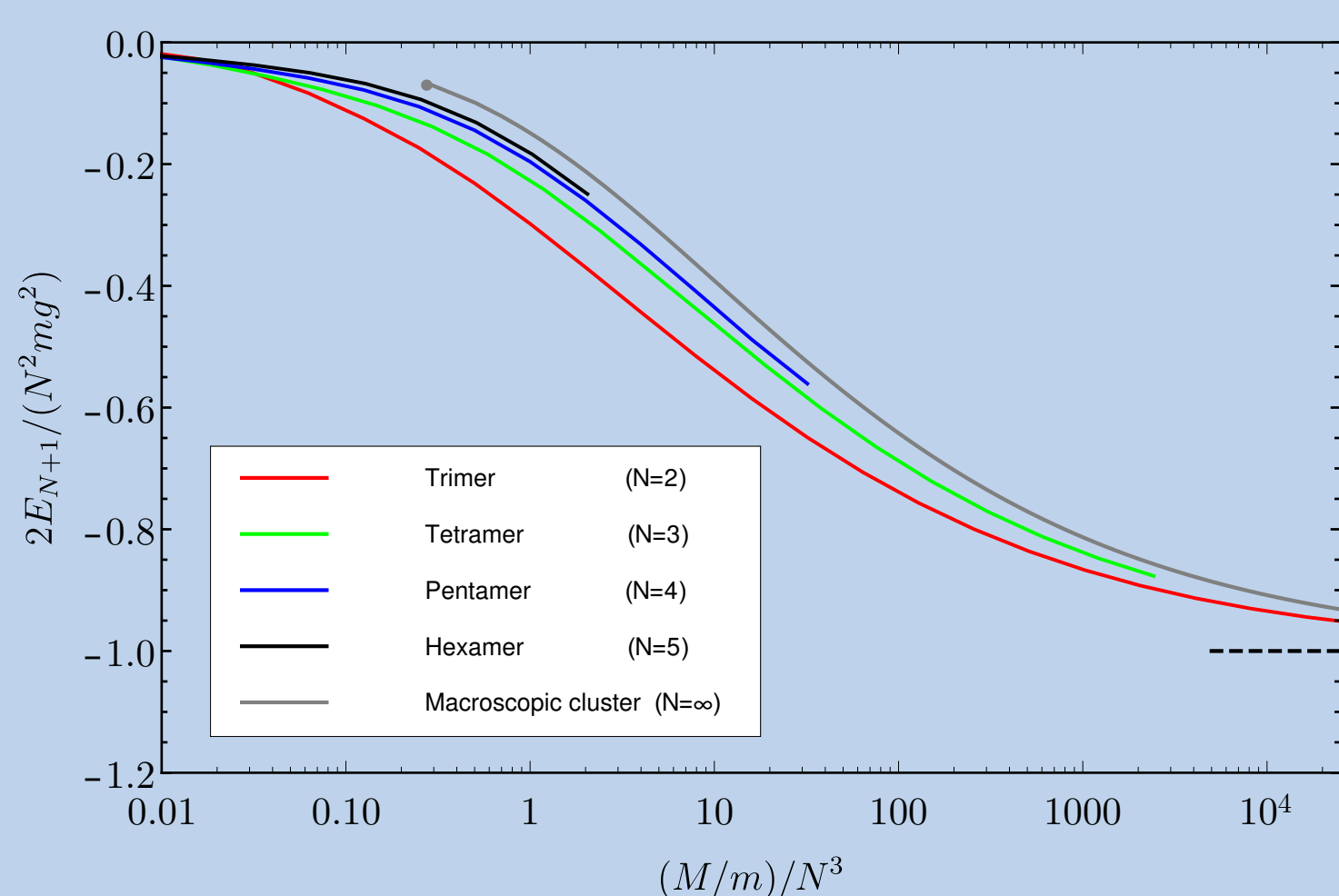
$$\Omega = \int \left[\frac{|\phi'(x)|^2}{2m} + gn(x)|\phi(x)|^2 + \frac{\pi^2 n^3(x)}{6M} - \epsilon |\phi(x)|^2 - \mu n(x) \right] dx$$

With normalization conditions

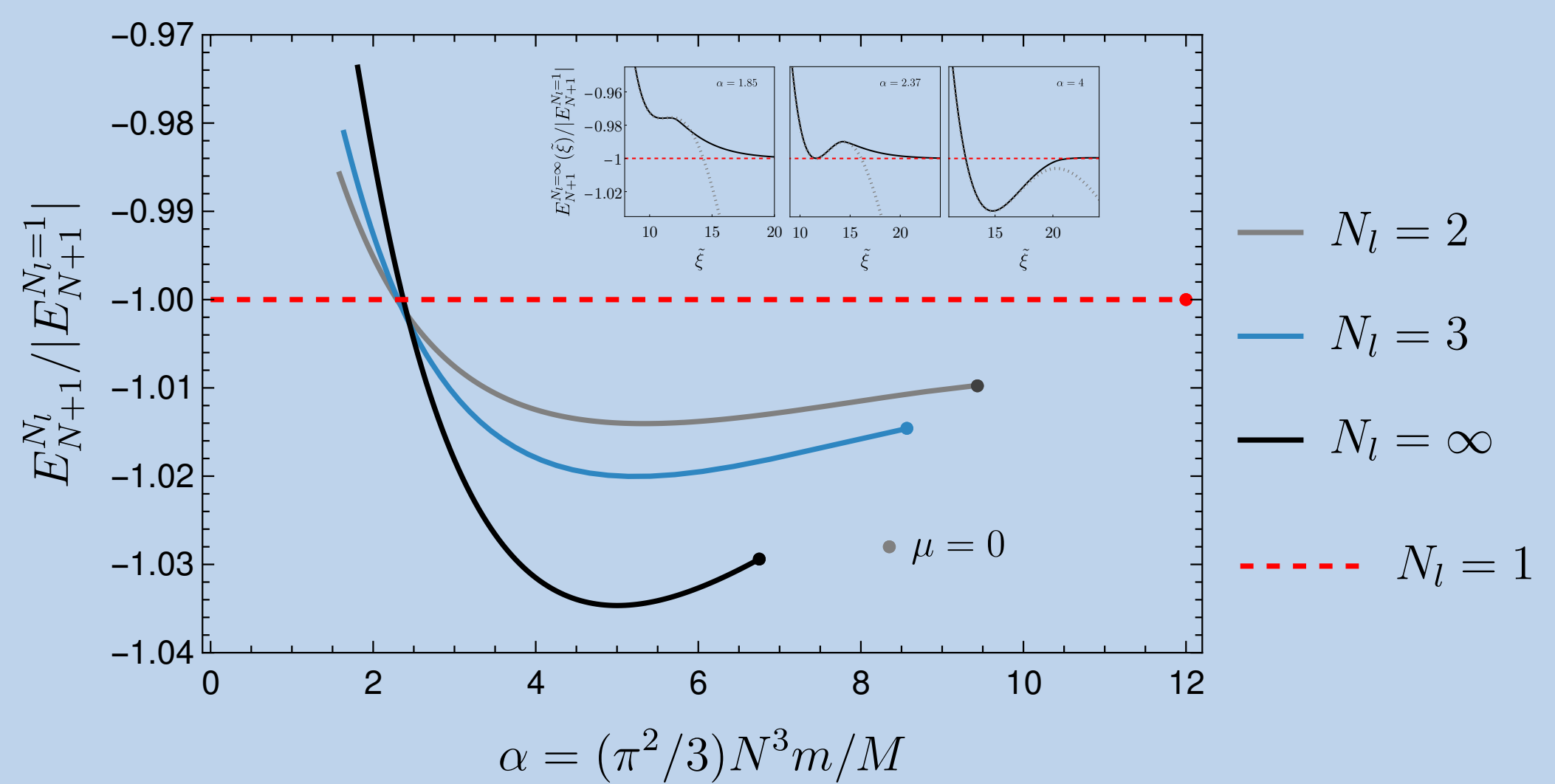
$$\begin{aligned} \int n(x) dx &= N \\ \int \phi^2(x) dx &= 1 \end{aligned}$$

When the chemical potential is zero, it can be solved analytically and we find:

$$\text{Critical mass-ratio for binding } (\mu=0): \\ (M/m)_{N+1} = \frac{\pi^2}{36} N^3$$



General case $N_h + N_l$ [2]



We follow the same mean-field approach with the density functional Thomas-Fermi approximation. We also introduce the lengthscale $\lambda = 1/(2m|g|N)$ with $N = N_h/N_l$. We use then use this length scale to rescale the different quantities:

$$\begin{aligned} \bullet \phi_i(x) &= \tilde{\phi}_i(u) & \bullet n(x) &= N \tilde{n}(u) & \bullet u &= x/\lambda \end{aligned}$$

Before	After
$\Omega = \int dx \left[\sum_{i=1}^{N_l} \left(\frac{ \partial_x \phi_i ^2}{2m} + gn \phi_i ^2 \right) + \frac{\pi^2 n^3}{6M} - \sum_{i=1}^{N_l} \epsilon_i \phi_i ^2 - \mu n \right]$	$\frac{\Omega}{2mg^2 N^2} = \int \left[\sum_{i=1}^{N_l} \left(\partial_u \tilde{\phi}_i ^2 - \tilde{n} \tilde{\phi}_i ^2 \right) + \alpha \tilde{n}^3 - \sum_{i=1}^{N_l} \tilde{\epsilon}_i \tilde{\phi}_i ^2 - \tilde{\mu} \tilde{n} \right] du$
<p>With normalization conditions</p> $\int n(x) dx = N_h$ $\int \phi_i(x) \phi_j(x) dx = \delta_{i,j}$	<p>With normalization conditions</p> $\int \tilde{n}(u) du = 1$ $\int \tilde{\phi}_i^*(u) \tilde{\phi}_j(u) du = \delta_{i,j}$

Single controlling parameter:

$$\alpha = \frac{\pi^2 m N^3}{3M}$$

We also consider the infinite chain $N_l \gg 1$ in a similar manner and solve everything numerically using an iterative scheme and find the following binding threshold for α :

Predicted range of α for binding

$$\begin{aligned} \alpha_{N_l=2} &\in [1.6, 9.4] \\ \alpha_{N_l=3} &\in [1.6, 8.6] \\ \alpha_{N_l=\infty} &\in [1.8, 6.8] \end{aligned}$$

References :

- [1] Binding of heavy fermions by a single light atom in one dimension, A. Tononi, J. Givois, D. S. Petrov, Phys. Rev. A 106, L011302 (2022)
[2] Self-binding of one-dimensional fermionic mixtures with zero-range interspecies attraction, J. Givois, A. Tononi, D. S. Petrov, arXiv:2207.04742